

# G-Function of Two Variables and the Conduction of Heat on a Flat Earth



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## Abstract

Various problems in science and technology, when formulated mathematically, lead naturally to certain classes of partial differential equations involving one or more unknown functions together with the prescribed conditions (known as boundary conditions) which arise from the physical situation. Several workers have obtained solutions to the equations related to certain problems, which satisfy the given boundary conditions. The classical method in obtaining solutions of the boundary value problems of mathematical physics can be derived from Fourier's. Another technique using integral transforms, which had its origin in Heaviside's work, has been developed in the past and has certain advantages over the classical method.

The object of this paper is to evaluate an integral involving G-function of two variables and employ it to derive a solution of the problem of the conduction of heat in a flat earth.

**Keywords:** Boundary Value Problems, G-Function of Two Variables, Conduction of Heat, Flat Earth.

## Introduction

The G-function of two variables was defined by Shrivastava and Joshi [4, p. 471] in terms of Mellin-Barnes type integrals as follows:

$$G_{p_1, q_1; p_2, q_2; p_3, q_3}^{0, n_1; m_2, n_2; m_3, n_3} (x, y) = \int_{L_1} \int_{L_2} \phi_1(\xi, \eta) \theta_2(\xi) \theta_3(\eta) x^\xi y^\eta d\xi d\eta \quad (1)$$

where

$$\phi_1(\xi, \eta) = \frac{\prod_{j=1}^{n_1} \Gamma(1-a_j+\xi+\eta)}{\prod_{j=n_1+1}^{p_1} \Gamma(a_j-\xi-\eta) \prod_{j=1}^{q_1} \Gamma(1-b_j+\xi+\eta)},$$

$$\theta_2(\xi) = \frac{\prod_{j=1}^{m_2} \Gamma(d_j-\xi) \prod_{j=1}^{n_2} \Gamma(1-c_j+\xi)}{\prod_{j=m_2+1}^{q_2} \Gamma(1-d_j+\xi) \prod_{j=n_2+1}^{p_2} \Gamma(c_j-\xi)}$$

$$\theta_3(\eta) = \frac{\prod_{j=1}^{m_3} \Gamma(f_j-\eta) \prod_{j=1}^{n_3} \Gamma(1-e_j+\eta)}{\prod_{j=m_3+1}^{q_3} \Gamma(1-f_j+\eta) \prod_{j=n_3+1}^{p_3} \Gamma(e_j-\eta)}$$

x and y are not equal to zero, and an empty product is interpreted as unity  $p_i, q_i, n_i$  and  $m_j$  are non negative integers such that  $p_i \geq n_i \geq 0, q_i \geq m_j \geq 0, (i = 1, 2, 3; j = 2, 3)$ .

The contour  $L_1$  is in the  $\xi$ -plane and runs from  $-i\infty$  to  $+i\infty$ , with loops, if necessary, to ensure that the poles of  $\Gamma(d_j - \xi)$  ( $j = 1, \dots, m_2$ ) lie to the right, and the poles of  $\Gamma(1 - c_j + \xi)$  ( $j = 1, \dots, n_2$ ),  $\Gamma(1 - a_j + \xi + \eta)$  ( $j = 1, \dots, n_1$ ) to the left of the contour.

The contour  $L_2$  is in the  $\eta$ -plane and runs from  $-i\infty$  to  $+i\infty$ , with loops, if necessary, to ensure that the poles of  $\Gamma(f_j - \eta)$  ( $j = 1, \dots, m_3$ ) lie to the right, and the poles of  $\Gamma(1 - e_j + \eta)$  ( $j = 1, \dots, n_3$ ),  $\Gamma(1 - a_j + \xi + \eta)$  ( $j = 1, \dots, n_1$ ) to the left of the contour.

and the double integral converges if

$$2(n_1 + m_2 + n_2) > (p_1 + q_1 + p_2 + q_2)$$

$$2(n_1 + m_3 + n_3) > (p_1 + q_1 + p_3 + q_3)$$

and  
where

$$| \arg x | < \frac{1}{2} U\pi, | \arg y | < \frac{1}{2} V\pi$$

$$U = [n_1 + m_2 + n_2 - \frac{1}{2} (p_1 + q_1 + p_2 + q_2)]$$

$$V = [n_1 + m_3 + n_3 - \frac{1}{2} (p_1 + q_1 + p_3 + q_3)]$$

These assumptions for the G-function of two variables will be adhered to throughout this research work.

The following formulae are required in the proof:  
Legendre's duplication formula [3]:

$$\sqrt{\pi}\Gamma(2z) = 2^{2z-1}\Gamma(z)\Gamma(z + \frac{1}{2}). \tag{2}$$

The following modified form of the integral [2]:

$$\int_{-T}^T e^{-2\pi i m t/T} \left(\cos \frac{\pi t}{2T}\right)^{\omega-1} dt = \frac{\Gamma(\omega)}{2^{\omega-2}\Gamma\left(\frac{\omega-4m+1}{2}\right)\Gamma\left(\frac{\omega+4m+1}{2}\right)}, \text{Re}(\omega) > -1. \tag{3}$$

**Aim of the Study**

To derive a solution of the problem of the conduction of heat in a flat earth involving G-function of two variables.

**Integral**

The integral, which we need as follows:

$$\int_{-T}^T e^{-2\pi i m t/T} \left(\cos \frac{\pi t}{2T}\right)^{\omega-1} G_{p_1, q_1; p_2, q_2; p_3, q_3}^{0, n_1; m_2, n_2; m_3, n_3} \left[ z_1 \left(\cos \frac{\pi t}{2T}\right)^2 \left| \begin{matrix} (a_j; 1, 1)_{1, p_1}; (c_j, 1)_{1, p_2}; (e_j, 1)_{1, p_3} \\ (b_j; 1, 1)_{1, q_1}; (d_j, 1)_{1, q_2}; (f_j, 1)_{1, q_3} \end{matrix} \right. \right] dt$$

$$= \frac{2T}{\sqrt{\pi}} G_{p_1, q_1; p_2+2, q_2+2; p_3, q_3}^{0, n_1; m_2, n_2+2; m_3, n_3} \left[ z_1 \left| \begin{matrix} (a_j; 1, 1)_{1, p_1}; \left(1-\frac{\omega}{2}, 1\right), \left(\frac{1-\omega}{2}, 1\right); (c_j, 1)_{1, p_2}; (e_j, 1)_{1, p_3} \\ (b_j; 1, 1)_{1, q_1}; (d_j, 1)_{1, q_2}; \left(\frac{1-\omega+4m}{2}, 1\right); (f_j, 1)_{1, q_3} \end{matrix} \right. \right], \tag{4}$$

provided that

$$\omega > 0, |\arg x| < \frac{1}{2} U\pi, |\arg y| < \frac{1}{2} V\pi$$

where

$$U = [n_1 + m_2 + n_2 - \frac{1}{2} (p_1 + q_1 + p_2 + q_2)]$$

$$V = [n_1 + m_3 + n_3 - \frac{1}{2} (p_1 + q_1 + p_3 + q_3)]$$

**Proof**

The integral (4) can be obtained easily by making use of the definition of G-function of two variables as given in (1) and the formulae (2) and (3).

**Formulate the Problem**

The surface of the earth is considered as a plane and f(t) is assumed as an averaged purely periodic temperature (annual averaged or daily averaged temperature). In order to determine the temperature in the interior of the earth ignoring the increase of temperature in the interior of the earth due to radio active or nuclear processes. We can, in general, use the method given in Sommerfeld [6] by setting  $x_0 = 0$  (surface of the earth),  $x_1 = \infty$  (great depth) and  $u_0 = f(t)$  for  $x = 0$ . It is convenient in this case to expand f(t) into a complex Fourier series:  $f(t) = \sum_{n=-\infty}^{\infty} C_n e^{2\pi i n t/T}$ , (T = length of year or day)

and to set for temperature in the interior of the earth at a depth x:

$$u(x, t) = \sum_{n=-\infty}^{\infty} C_n u_n(x) e^{2\pi i n t/T} \tag{6}$$

Each individual term of the series must satisfy the basic equation of heat conduction, viz.,  $\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2}$ . Thus yields the ordinary differential equation for  $u_n$ .

$$\frac{d^2 u_n}{dx^2} = p_n^2 u_n \text{ where } p_n^2 = \frac{2\pi i n}{kT} \tag{7}$$

In order for (6) to go into (5) for  $x = 0$ , we must have  $u_n(0) = 1$ .

Depending on whether n is positive or negative, we take

$$2in = (1 + i)^2 |n| \tag{8}$$

and

$$p_n = (1 + i)q_n, q_n = \sqrt{\frac{|n|\pi}{kT}} > 0. \tag{9}$$

The general solution of (7) is obtained as:

$$u_n(x) = A_n e^{(1+i)q_n x} + B_n e^{-(1+i)q_n x}. \tag{10}$$

In equation (10), we must have  $A_n = 0$ , since otherwise the temperature would become infinite for n

**Methodology**

Beg [1], Shrivastava [5] and several other authors have obtained solutions of various boundary value problems involving G-Function.

Following Beg [1], Shrivastava [5] and several other authors, we will employ the G-function of two variables in obtaining a solution of a boundary value problem related to conduction of heat in a flat earth, which will be useful for further research.

$\rightarrow \infty$  and  $B_n = 1$ , to satisfy (8). Substituting this in (6), we obtain our solution as [4]:

$$u(x, t) = \sum_{n=-\infty}^{\infty} C_n e^{-(1+i)q_n x} e^{2\pi i n t/T}. \tag{11}$$

In this paper, we have

$$u(0, t) = f(t) =$$

$$\left(\cos \frac{\pi t}{2T}\right)^{\omega-1} G_{p_1, q_1; p_2, q_2; p_3, q_3}^{0, n_1; m_2, n_2; m_3, n_3} \left[ z_1 \left(\cos \frac{\pi t}{2T}\right)^2 \left| \begin{matrix} (a_j; 1, 1)_{1, p_1}; (c_j, 1)_{1, p_2}; (e_j, 1)_{1, p_3} \\ (b_j; 1, 1)_{1, q_1}; (d_j, 1)_{1, q_2}; (f_j, 1)_{1, q_3} \end{matrix} \right. \right] \tag{12}$$

**Solution of the Problem**

The solution of the problem to be obtained is

$$u(x, t) = \frac{1}{\sqrt{\pi}} \sum_{n=-\infty}^{\infty} e^{-(1+i)q_n x} e^{2\pi i n t/T} \times$$

$$G_{p_1, q_1; p_2+2, q_2+2; p_3, q_3}^{0, n_1; m_2, n_2+2; m_3, n_3} \left[ z_1 \left| \begin{matrix} (a_j; 1, 1)_{1, p_1}; \left(1-\frac{\omega}{2}, 1\right), \left(\frac{1-\omega}{2}, 1\right); (c_j, 1)_{1, p_2}; (e_j, 1)_{1, p_3} \\ (b_j; 1, 1)_{1, q_1}; (d_j, 1)_{1, q_2}; \left(\frac{1-\omega+4n}{2}, 1\right); (f_j, 1)_{1, q_3} \end{matrix} \right. \right] \tag{13}$$

where all conditions of convergence are same as in (4).

**Proof**

If  $t = 0$ , then by virtue of (11) and (12), we have

$$\sum_{n=-\infty}^{\infty} C_n e^{2\pi i n t/T} =$$

$$\left(\cos \frac{\pi t}{2T}\right)^{\omega-1} G_{p_1, q_1; p_2, q_2; p_3, q_3}^{0, n_1; m_2, n_2; m_3, n_3} \left[ z_1 \left(\cos \frac{\pi t}{2T}\right)^2 \left| \begin{matrix} (a_j; 1, 1)_{1, p_1}; (c_j, 1)_{1, p_2}; (e_j, 1)_{1, p_3} \\ (b_j; 1, 1)_{1, q_1}; (d_j, 1)_{1, q_2}; (f_j, 1)_{1, q_3} \end{matrix} \right. \right] \tag{14}$$

Multiplying both side of (14) by  $e^{-2\pi i m t/T}$  and integrating with respect to t from  $-T$  to  $T$ , we get

$$\int_{-T}^T e^{-2\pi i m t/T} \left(\cos \frac{\pi t}{2T}\right)^{\omega-1} \times G_{p_1, q_1; p_2, q_2; p_3, q_3}^{0, n_1; m_2, n_2; m_3, n_3} \left[ z_1 \left(\cos \frac{\pi t}{2T}\right)^2 \left| \begin{matrix} (a_j; 1, 1)_{1, p_1}; (c_j, 1)_{1, p_2}; (e_j, 1)_{1, p_3} \\ (b_j; 1, 1)_{1, q_1}; (d_j, 1)_{1, q_2}; (f_j, 1)_{1, q_3} \end{matrix} \right. \right] dt$$

$$= \sum_{n=-\infty}^{\infty} C_n \int_{-T}^T e^{2\pi i(n-m)t/T} dt.$$

Now using (4) and orthogonality property of the exponential functions, viz.

$$\int_{-T}^T e^{2\pi i(n-m)t/T} dt = \begin{cases} 0, & \text{if } m \neq n; \\ 2T, & \text{if } m = n; \end{cases}$$

we get

$$C_m =$$

$$\frac{1}{\sqrt{\pi}} G_{p_1, q_1; p_2+2, q_2+2; p_3, q_3}^{0, n_1; m_2, n_2+2; m_3, n_3} \left[ \begin{matrix} (a_j; 1, 1)_{1, p_1} : (1 - \frac{\omega}{2}, 1), (\frac{1-\omega}{2}, 1), (c_j, 1)_{1, p_2} : (e_j, 1)_{1, p_3} \\ (b_j; 1, 1)_{1, q_1} : (d_j, 1)_{1, q_2}, (\frac{1-\omega+4m}{2}, 1) : (f_j, 1)_{1, q_3} \end{matrix} \right] \quad (15)$$

With the help of (11) and (15), the solution (13) is obtained.

**Conclusion**

Since the G-functions play a crucial role in a certain useful mathematical enterprise. When looked at conceptually, they are both natural and attractive. Most special functions, and many products of special functions, are G-functions or are expressible as products of G-functions with elementary functions.

On specializing the parameters, G-function of two variables may be reduced to several other higher transcendental functions. Therefore the result (13) is of general nature and may reduce to be in

different forms, which will be useful in the literature on applied Mathematics and other branches.

**References**

1. Beg Mehphooj and Shrivastava, S. S.: *Multivariable G-Function and Flux Condition, Aryabhatta Journal of Mathematics and Informatics, Vol.08 Issue-02, (July - December, 2016), pp. 96-99.*
2. Gradshteyn, I. S. & Ryzhik's, I. M.: *Table of Integrals, Series and Products, Academic Press, New York (1980), 372.*
3. Rainville, E. D.: *Special Functions, Macmillan, New York, 1960.*
4. Shrivastava, H. M. & Joshi, C. M.: *Integration of Certain Products Associated with a Generalized Meijer Function, Proc. Cambridge Philos. Soc. 65, 471-477.*
5. Shrivastava, S. S. and Sikarwar Pinki: *Application of G-Function to solve the problem related to Cooling of a Sphere, International Journal of Advanced Technology & Engineering Research (IJATER), 2017, pp.57-59.*
6. Sommerfield, A.: *Partial Differential Equations in Physics, Academic Press, Inc., New York (1949).*